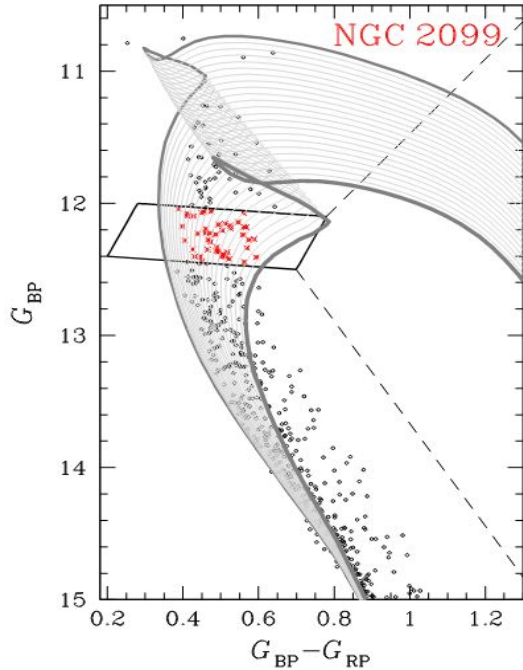

Code & stellar physics
workshop:

Study of intensity spectrum
from rapidly rotating stars

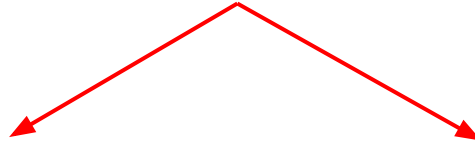
Lazzarotto Axel
PhD student , *IRAP*. 14, Avenue Edouard
Belin 31400 Toulouse.
Mail: alazzarotto@irap.omp.eu

State of the art

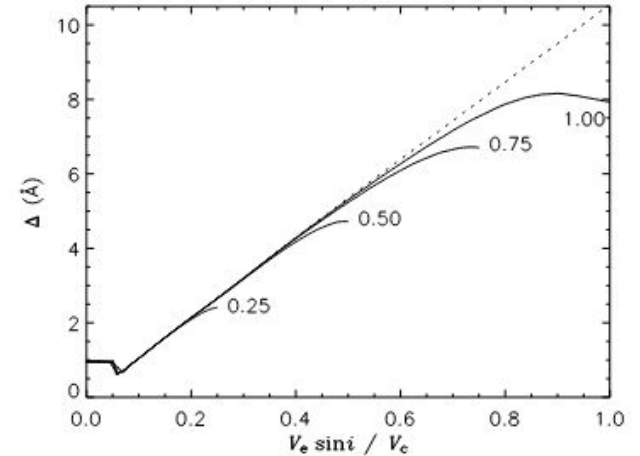
G. Cordoni et al. (2018)
Broadening of the isochron in the
vicinity of the turn-of



1D models

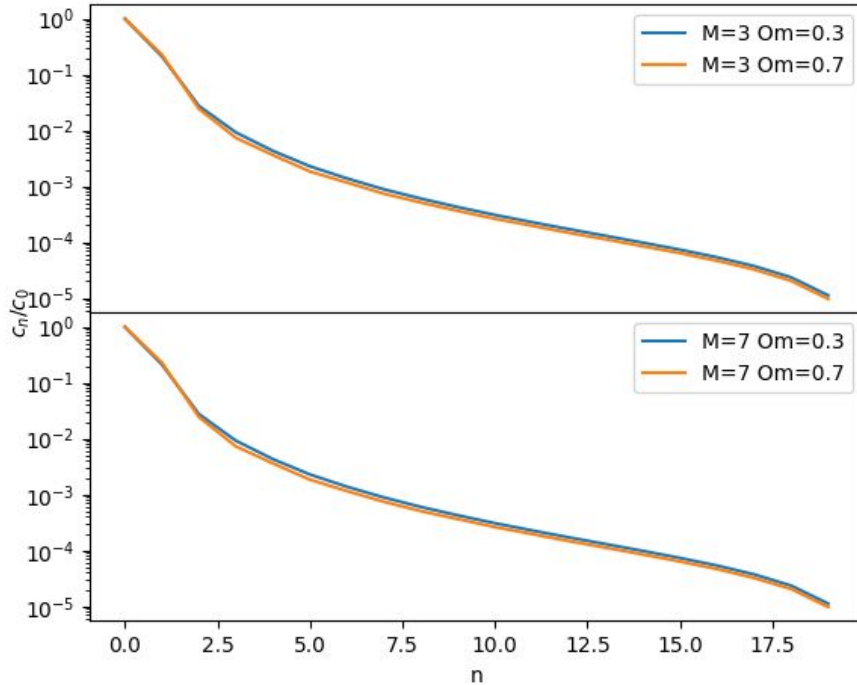


Townsend (2004) the equatorial
projected velocity is
underestimated



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ESTER : robustness of the latitudinal profiles



Interpolation of $(T_{\text{eff}}, \log g, R)$ by Legendre polynomials
=> Polynomials degree correspond to latitudinal resolution of ESTER $n\theta$

Choose $n\theta$ such as the interpolated T_{eff} ,
 $\forall \theta \mid |T_{\text{eff}}(\theta, n\theta) - T_{\text{eff}}(\theta, 20)| / T_{\text{eff}}(\theta, 20) < 1$

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PHOENIX

=> Solves Radiative Transfer Equation (RTE) under Local Thermodynamic Equilibrium (LTE) and non-LTE (NLTE) assumption(s).

=> Solves Saha-Boltzmann, charge conservation and statistic equilibrium equations

Assumptions: - Spherical symmetry (Matching with plane-parallel when atmosphere thickness \ll radius)

- Stationary state

=> **Pros:** - Give access to $I(\lambda, \mu)$ taking Limb-Darkening (LD) into account.

=> **Cons:** - Don't take Gravity Darkening (GD) into account

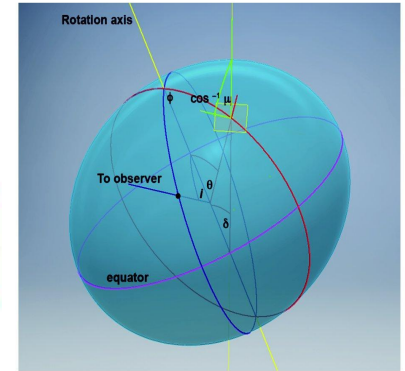
plane-parallel geometry:

$$\frac{dI_\nu}{ds} = -\mu \frac{dI_\nu}{dt} \quad \text{with } dt = -\mu$$

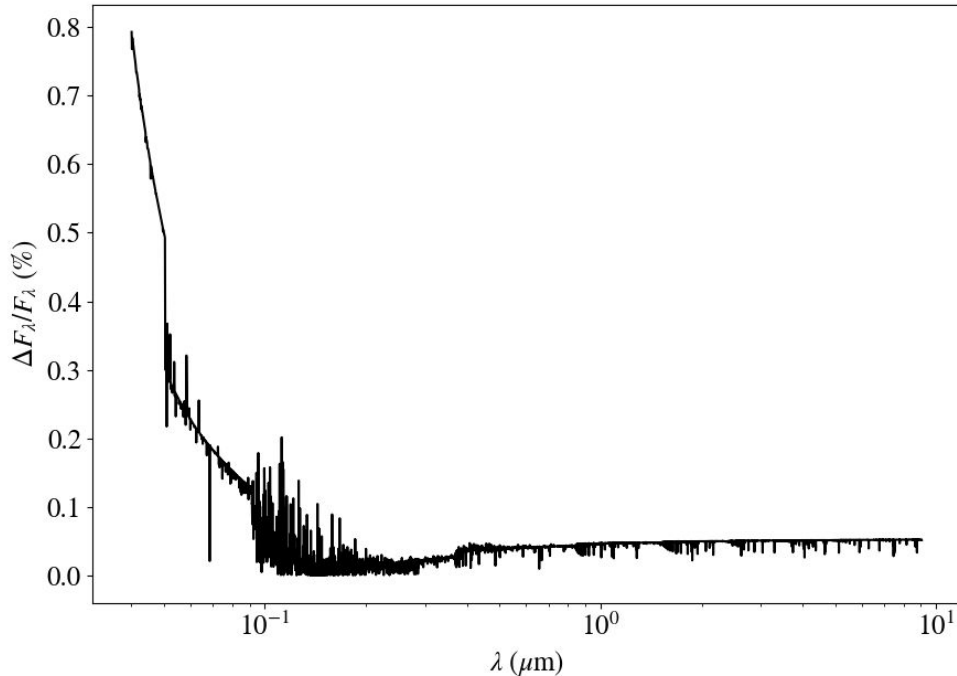
spherical geometry:

$$\frac{dI_\nu}{ds} = \frac{\partial I_\nu}{\partial r} \frac{dr}{ds} + \frac{\partial I_\nu}{\partial \mu} \frac{d\mu}{ds}$$

$$\Rightarrow \frac{dI_\nu}{ds} = \frac{\partial I_\nu}{\partial r} \mu + \frac{\partial I_\nu}{\partial \mu} \frac{1-\mu^2}{r}$$

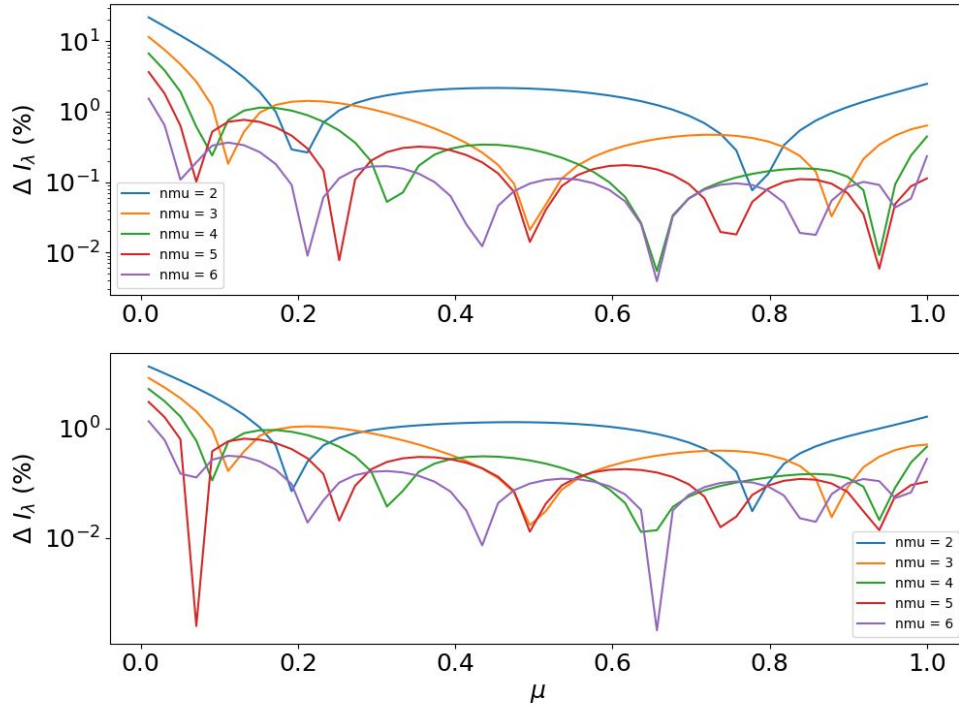


PHOENIX :
Parallel plan or spherical
model ?



The relative radiative flux difference between a plane-parallel and a spherical model is below 0.1 % in the wavelength interval we use. We, thus, have done all the calculations with the plane-parallel geometry to reduce the computing time

PHOENIX : Gauss grid



- $N_\mu = 4$ ensure a numerical precision better than 5% everywhere (less than 1% mostly)

Rebuilding spectrum:
Spatial grid

- Surface element: $\Delta S(\theta_k) = r^2(\theta_k) \sqrt{1 + \frac{r_\theta^2(\theta_k)}{r^2(\theta_k)}} \sin \theta_k \Delta\theta \Delta\phi$

$$r_\theta(\theta_k) = \partial r / \partial \theta$$
$$\Delta\phi = 2\pi / N_\phi(\theta_k)$$

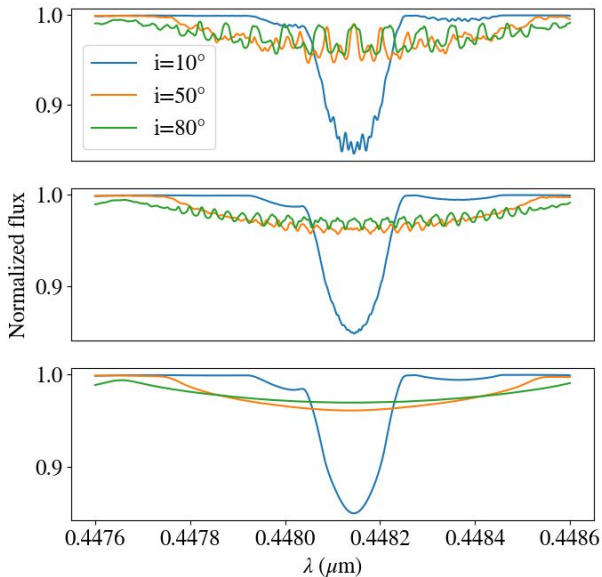
=> All elements must have the same ΔS in order to avoid an overestimation of polar contributions and an underestimation of the equatorial ones.

=> We thus define our grid with N_θ and $N_\phi(\Theta_0)$

- Visible part of the grid : $\mu \geq 0$

$$\mu(\theta_k, \phi_j) = \frac{1}{\sqrt{r^2 + r_\theta^2}} \left[(r \sin \theta_k - r_\theta \cos \theta_k) \cos \phi_j \sin i + (r \cos \theta_k + r_\theta \sin \theta_k) \cos i \right]$$

Rebuilding spectrum: Spatial grid



$$(N_{\Theta}, N_{\phi}(\Theta_0)) = (10, 10)$$

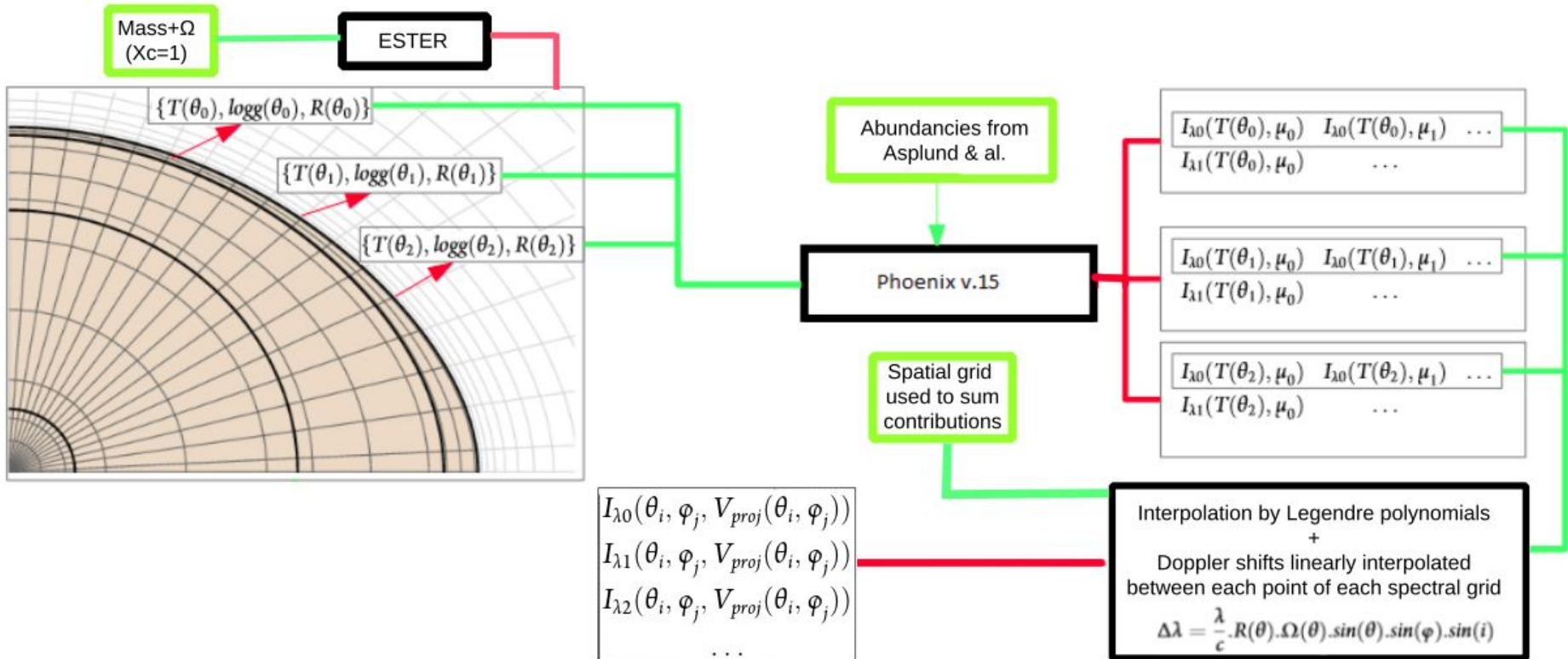
$$(N_{\Theta}, N_{\phi}(\Theta_0)) = (20, 10)$$

$$(N_{\Theta}, N_{\phi}(\Theta_0)) = (100, 20)$$

- Oscillations due to a lack of spatial resolution which lead to a bad mapping of Doppler effect.

=> Optimal N_{Θ} and $N_{\phi}(\Theta_0)$ depend on mass, rotation and inclination

Rebuilding spectrum:

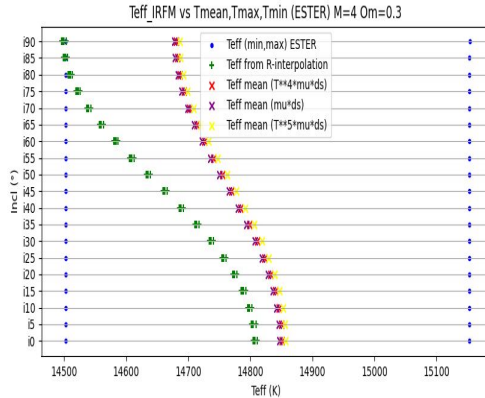


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Possible applications

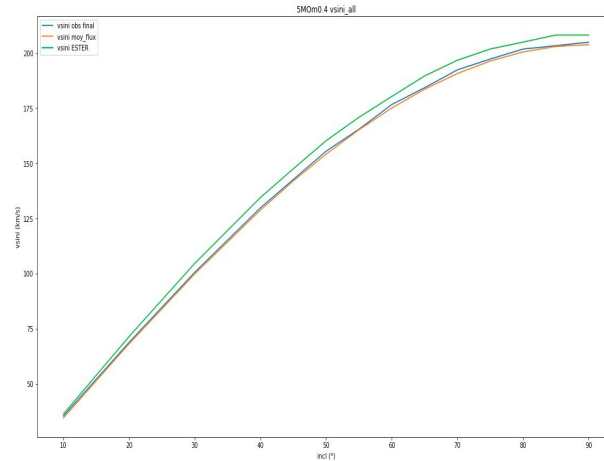
- Photometry:

Example : Study the variation of an observable according the inclination.



- Spectrometry:

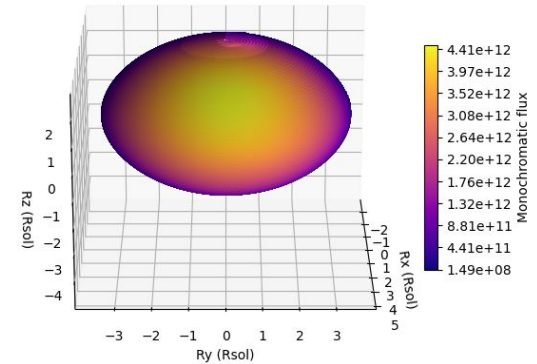
Example : Study of a measured equatorial velocity



- Interferometry:

Example : Study of monochromatic maps of the st

Monochromatic ($\lambda=3595\text{\AA}$) star ($\text{Om}=6, i=0.7$) viewed from Earth



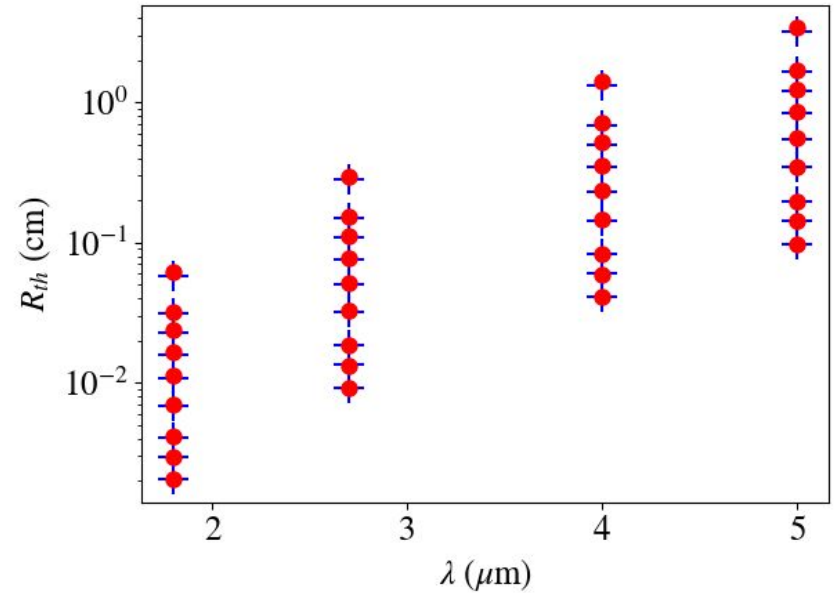
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InfraRed Method Flux: IRFM

Method based on the study
of the following ratio:

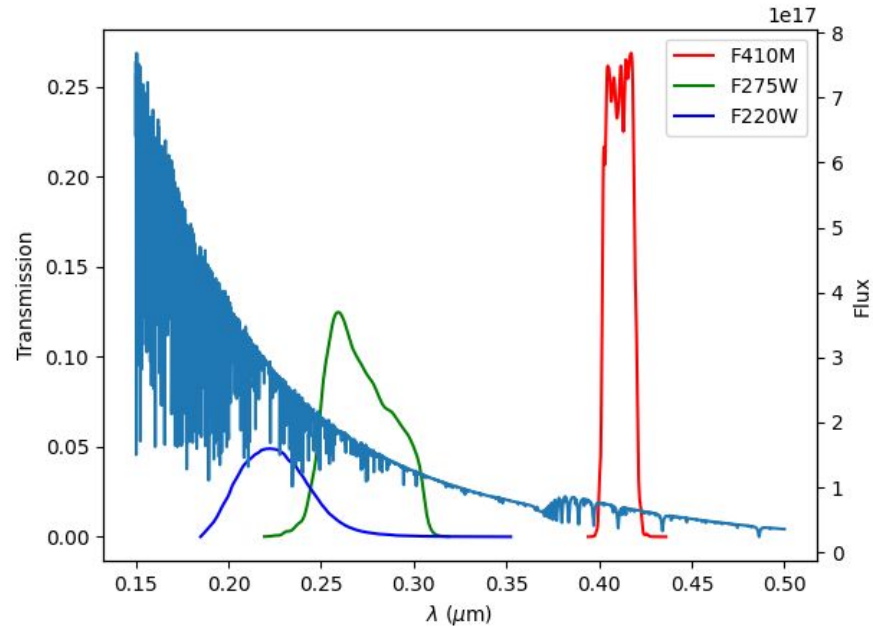
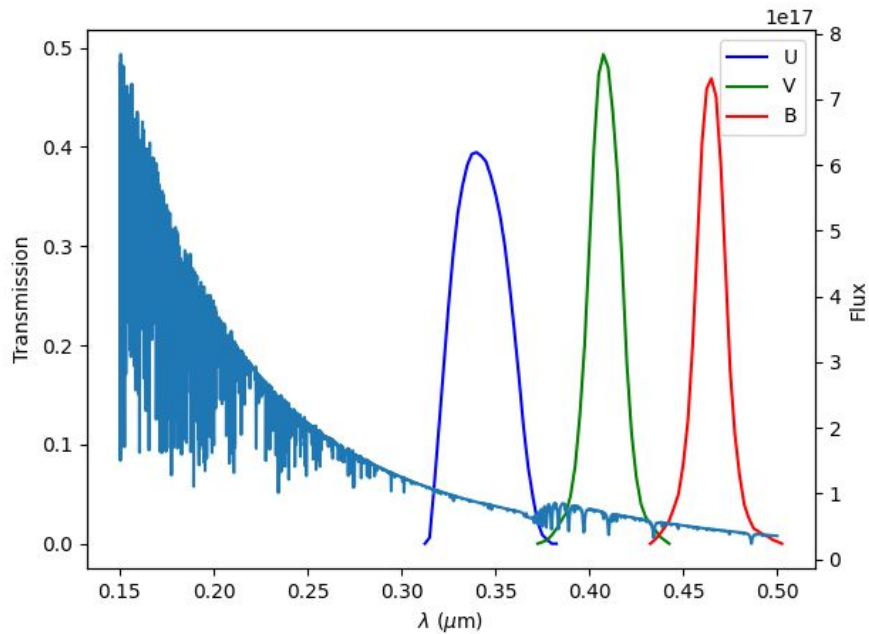
$$R = \frac{\int_0^{\infty} F_{\lambda} d\lambda}{F_{IR}}$$

with an IR wavelength $> 2\mu\text{m}$ to be
in the Rayleigh domain.



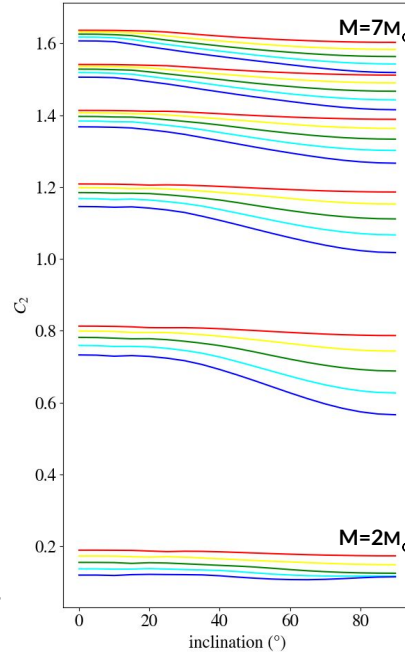
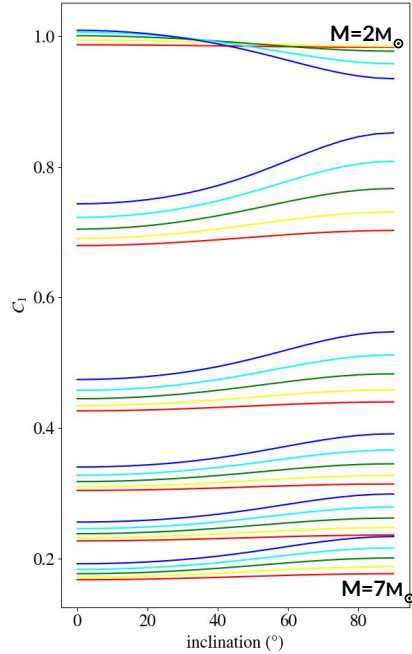
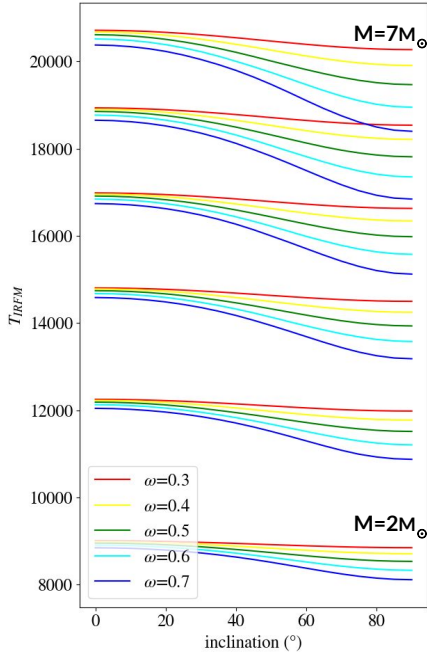
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Color indexes



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Tricubic interpolation



$$T_{IRFM}(M, \Omega, i) = \sum_{k=1}^{n_M} \sum_{l=1}^{n_{\omega}} \sum_{m=1}^{n_{\sin i}} a_{k,l,m}^{IRFM} B_{k,l,m}^{IRFM}(T_{sph}, \omega, \sin i)$$

$$c_1(M, \Omega, i) = \sum_{k=1}^{n_M} \sum_{l=1}^{n_{\omega}} \sum_{m=1}^{n_{\sin i}} a_{k,l,m}^{c_1} B_{k,l,m}^{c_1}(T_{sph}, \omega, \sin i)$$

$$c_2(M, \Omega, i) = \sum_{k=1}^{n_M} \sum_{l=1}^{n_{\omega}} \sum_{m=1}^{n_{\sin i}} a_{k,l,m}^{c_2} B_{k,l,m}^{c_2}(T_{sph}, \omega, \sin i)$$

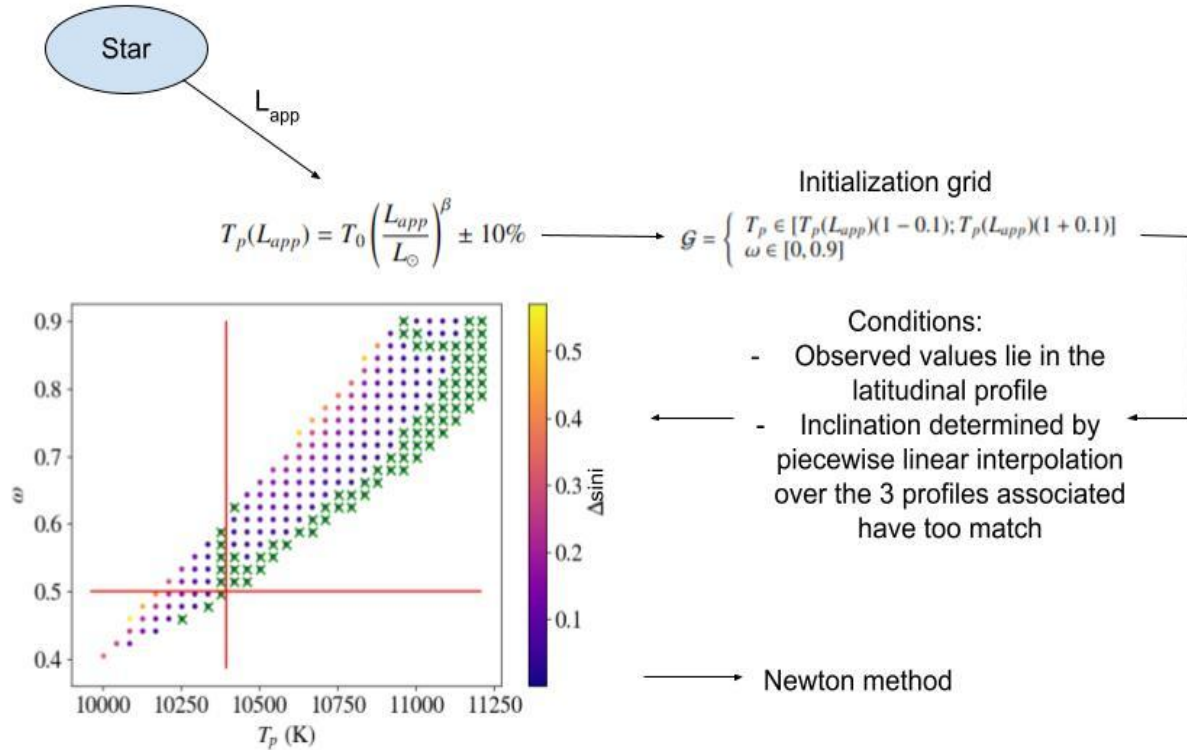
=> Interpolation by Tricubic splines

$$T_{sph} = T_p \left(1 - \sum_{k=1}^{n_M} a_k B_k(\omega) \right)$$

=> Orthogonalization of the parameter space

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Initialisation :



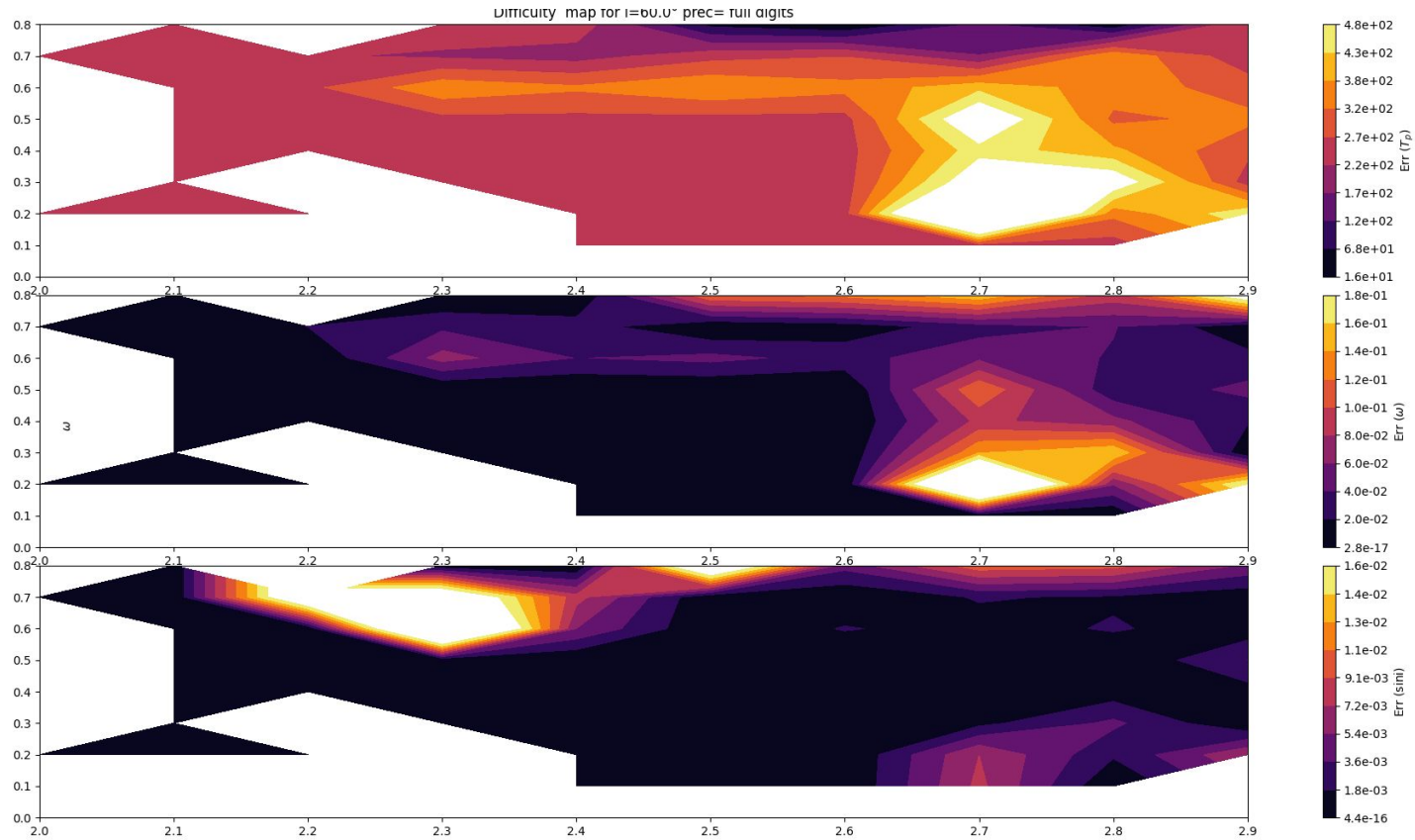
Resolution : feasibility map



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15

Resolution : error map



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Application to Vega

Grid of models between $M=[2 ; 2.9]M_{\odot}$ $\omega=[0 ; 0.8]$ $X_c = 0.268$

Vega model :

- $M = 2.36 M_{\odot}$

- $\omega = 0.51$

- $i = 6^{\circ}$

Observable from Vega :

- $T_{\text{IRFM}} = 9553 \text{ K}$ (Ciardi & al 2001)

- $C_1 = 1.089$ (Crawford & Barnes 1970)

- $L_{\text{app}} = 40.12L_{\odot}$ (Yoon & al 2010)

- $C_2 = 0.681$ (Compute from the concordance model from Lara & Rieutord 2013)

Solution:

T_p : 10073 K

Omega : 0.48

sini : 13°

Thanks for your attention !