

EWASS 2016, Symposium 17 :

Magnetic Helicity in Sun and Stars: From Dynamo Action to Eruptive Phenomena

Athens, Greece, July 8th 2016

Magnetic helicity conservation in a solar active event

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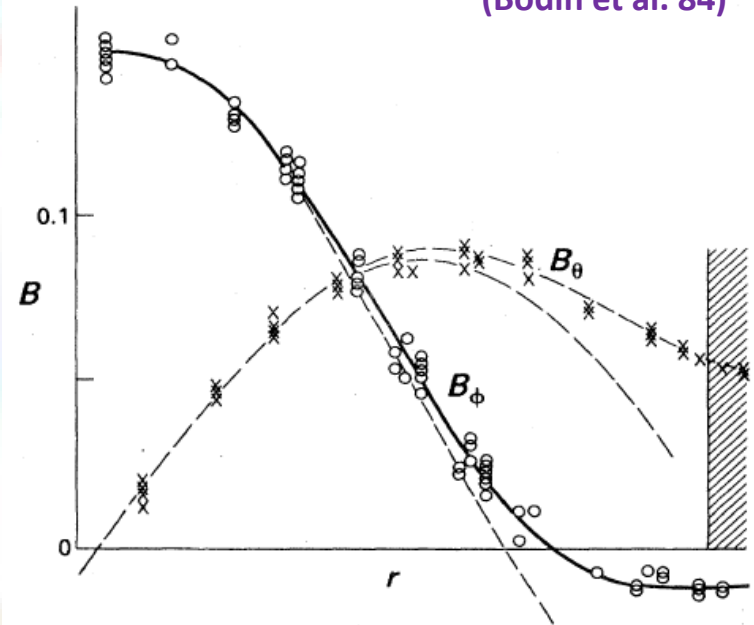
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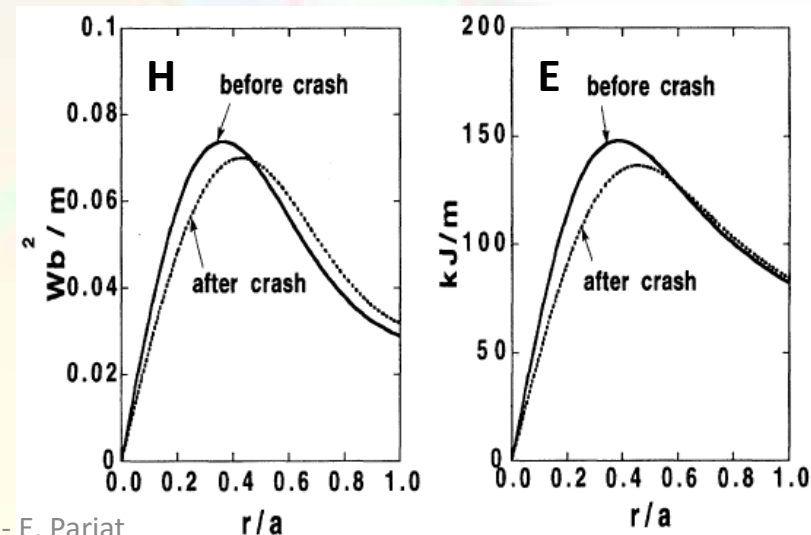
Taylor conjecture

- Magnetic helicity is an ideal MHD invariant when field lines cannot reconnect. For $E \perp B$: no dissipation \rightarrow magnetic helicity is conserved (Woltjer 58).

(Bodin et al. 84)



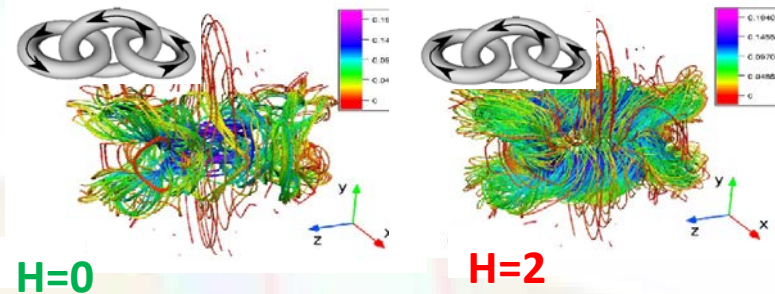
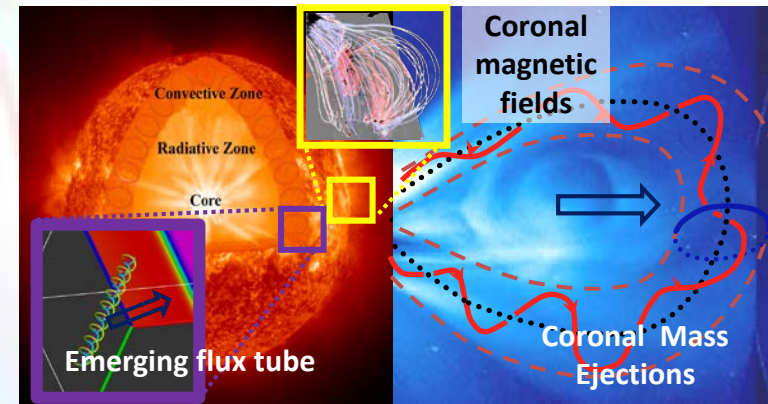
(Ji et al. 95)



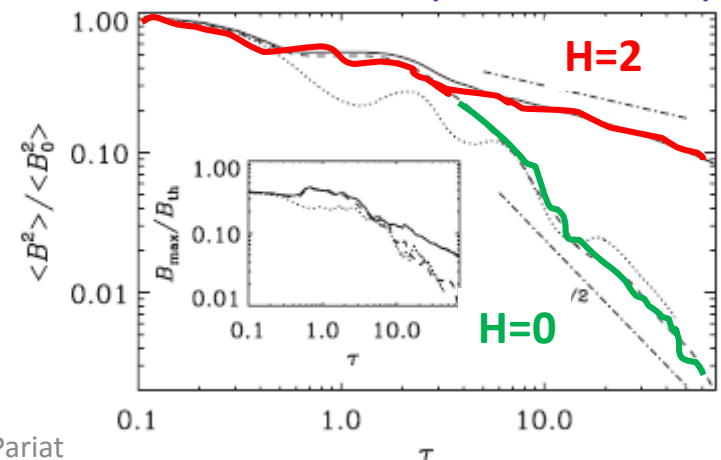
- Time variations** **Surface Flux** **Dissipation**
- $$\frac{dH_m}{dt} = \int_{\partial V} \left(\mathbf{A} \times \frac{\partial \mathbf{A}}{\partial t} \right) \cdot d\mathbf{S} - 2 \int_{\partial V} (\mathbf{E} \times \mathbf{A}) \cdot d\mathbf{S} - 2 \int_V \mathbf{E} \cdot \mathbf{B} \, dV$$
- From laboratory plasma experiment, **Taylor (1974) conjectured: even in non-ideal MHD magnetic helicity should be well conserved.**
 - Magnetic energy cascades to small scales where it is dissipated vs helicity cascades to large scales (Ji et al. 95, Heidbrink & Dang 00).
 - In resistive MHD, helicity dissipation is bounded and slow compared to energy dissipation (Berger 84, Berger 99)

Consequence of Taylor's conjecture

- Relaxation in lab. experiments: plasma relax to minimum energy state, i.e. linear force free field (LFFF) e.g. Bodin et al. 84, Taylor et al. 86, Yamada et al. 99
- Impact on dynamic of magnetic reconnection: e.g. Linton et al. 01, Del Soro et al. 10
- Helicity conservation plays role in:
 - Saturation of stellar dynamos: e.g. Brandenburg et al. 05, Blackman 14
 - Coronal heating models: Heyvaerts & Priest 84
 - Coronal mass ejection generation
- **Magnetic helicity conservation is the “raison d'être” of CMEs (Rust 94, Low 96)**
 - Conjecture : to limit the buildup of magnetic helicity in the corona, magnetic helicity has to be ejected via CMEs.

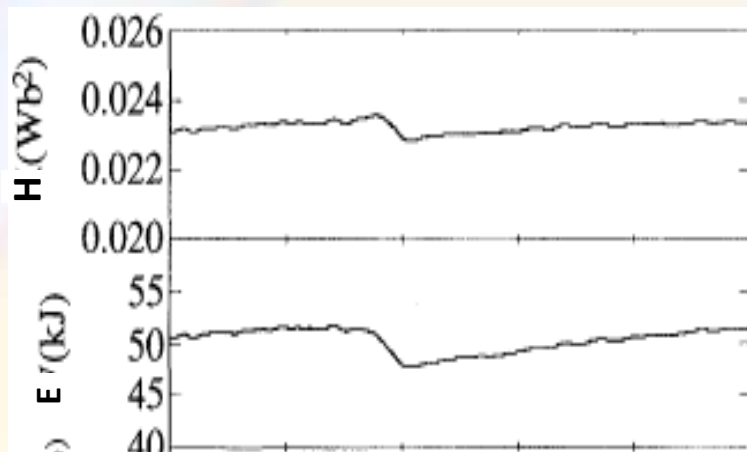


(Del Soro et al. 10)

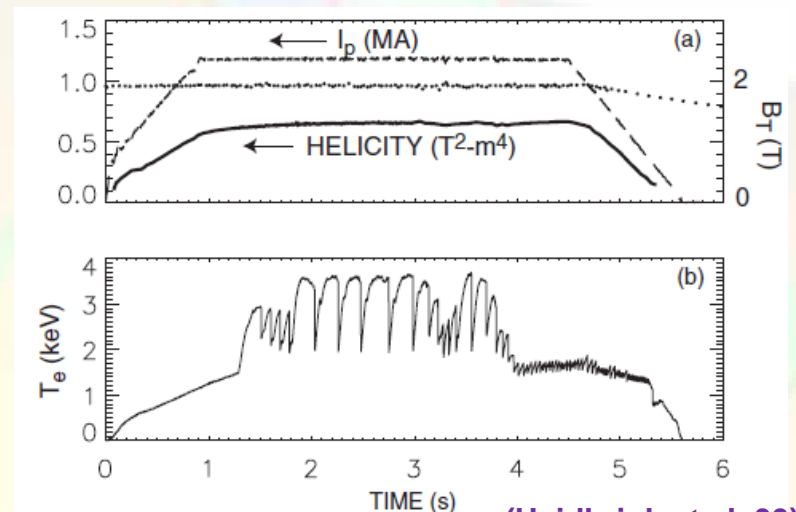


Tests on Magnetic Helicity Conservation

- **Despite its potential importance, tests on Taylor's conjecture have been very limited!**
 - Test on "relaxation" toward minimum energy state (LFFF): mixed results
 - ➔ not direct test of magnetic helicity conservation, but of relaxation dynamics
 - Laboratory experiments: difficult sampling of the full 3D magnetic field ; axisymmetric assumption (Ji et al. 95, Barnes et al. 86, Heidbrink et al. 00, Gray et al. 10)
 - Sawtooth relaxation: $\Delta H/H=1-5\%$; $\Delta E/E=5-10\%$
 - Sawtooth crash: $\Delta H/H=1\%$
- **Numerical simulation: no test in general conditions**, i.e. in 3D, active-like conditions, no periodicity ...



(Ji et al. 95)



(Heidbrink et al. 00)

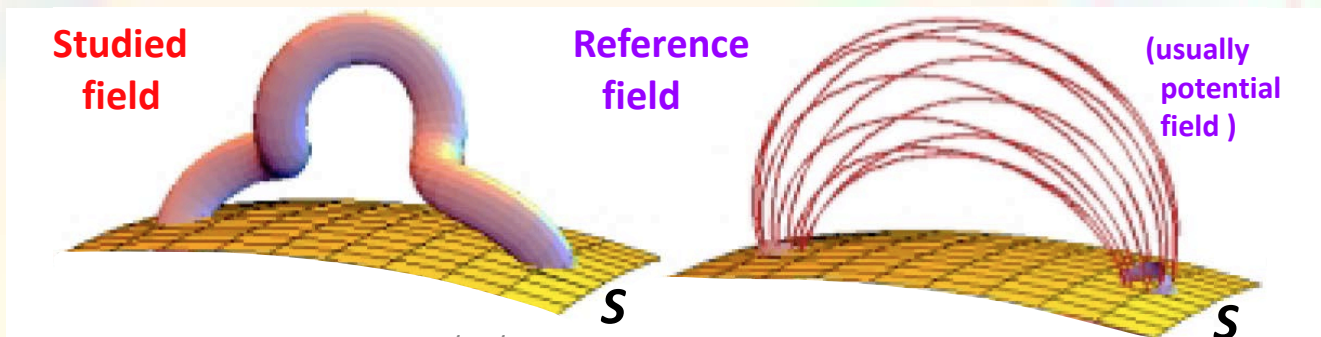
Relative Magnetic Helicity

- Magnetic helicity is gauge invariant only for magnetically bounded systems:
 - Strict definition of magnetic helicity useless for numerous applications:
 - e.g. natural plasmas, like the solar corona have boundaries threaded by magnetic fields
- Useful quantity: **Relative Magnetic Helicity**: helicity of a studied field relative to a reference field (Berger 84).

$$H_{\mathcal{V}} = \int_{\mathcal{V}} (\mathbf{A} + \mathbf{A}_p) \cdot (\mathbf{B} - \mathbf{B}_p) d\mathcal{V} \quad (\text{Finn \& Antonsen 85})$$

with boundary condition : $(\mathbf{B}_p \cdot d\mathbf{S})|_{\partial\mathcal{V}} = (\mathbf{B} \cdot d\mathbf{S})|_{\partial\mathcal{V}} \quad \nabla \times \mathbf{A} = \mathbf{B}$

- **Gauge invariant provided that studied and reference fields share the same magnetic-flux distribution on the boundary.**



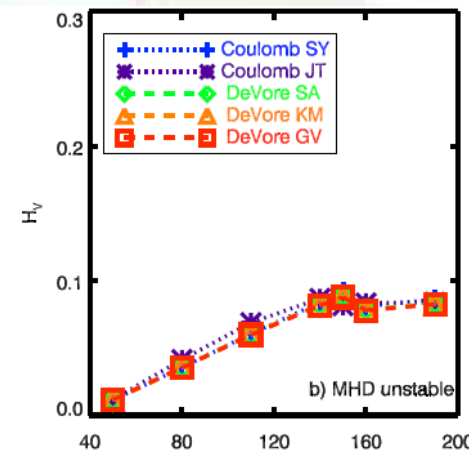
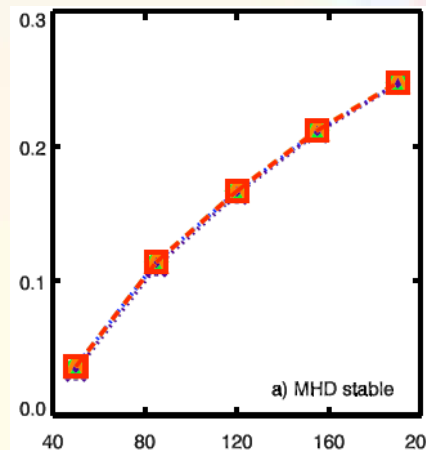
Relative Magnetic Helicity Estimations

- The computation of relative magnetic helicity on a 3D cuboid system has not been straightforward
 - Volume computation of consistent gauges for the studied and reference \mathbf{B}
 - Impose boundary conditions simultaneously on all 6 faces
- Methods recently developed:
 - Using Coulomb gauge: $\nabla \cdot \mathbf{A} = 0$
 - Thalmann et al. 11, Rudenko & Myshyakov 11, Yang et al. 13
 - Simpler theoretical formulation but harder to implement numerically
 - Using DeVore gauge: $A_z = 0$
 - Valori, Démoulin & Pariat 12
 - More complex theoretical formulation but simpler numerical implement.

• **Benchmarking of these methods performed by ISSI team on "Helicity estimations in models and observations"**

→ Cf. Talk of G. Valori

• **Methods perform very consistently when \mathbf{B} sufficiently solenoidal**



Magnetic helicity dissipation estimation

- General formulation of the time variation of the relative magnetic helicity:

Magnetic helicity dissipation

Time
variation
of relative
magnetic
helicity

$$\frac{dH}{dt} = -2 \int_{\mathcal{V}} \mathbf{B} \cdot \mathbf{E} \, d\mathcal{V} + 2 \int_{\mathcal{V}} \frac{\partial \phi}{\partial t} \nabla \cdot \mathbf{A}_p \, d\mathcal{V} + \int_{\partial \mathcal{V}} \left((\mathbf{A} - \mathbf{A}_p) \times \frac{\partial (\mathbf{A} + \mathbf{A}_p)}{\partial t} \right) \cdot d\mathbf{S} - 2 \int_{\partial \mathcal{V}} \frac{\partial \phi}{\partial t} \mathbf{A}_p \cdot d\mathbf{S} + 2 \int_{\partial \mathcal{V}} (\mathbf{B} \cdot \mathbf{A}_p) \mathbf{v} \cdot d\mathbf{S} - 2 \int_{\partial \mathcal{V}} (\mathbf{v} \cdot \mathbf{A}_p) \mathbf{B} \cdot d\mathbf{S}$$

Helicity variation and flux
of the reference field

Flux of helicity of the studied field

(Pariat et al. 15)

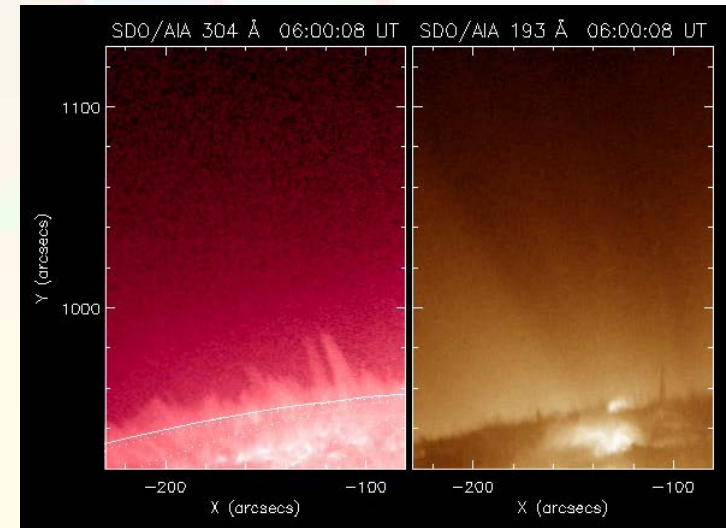
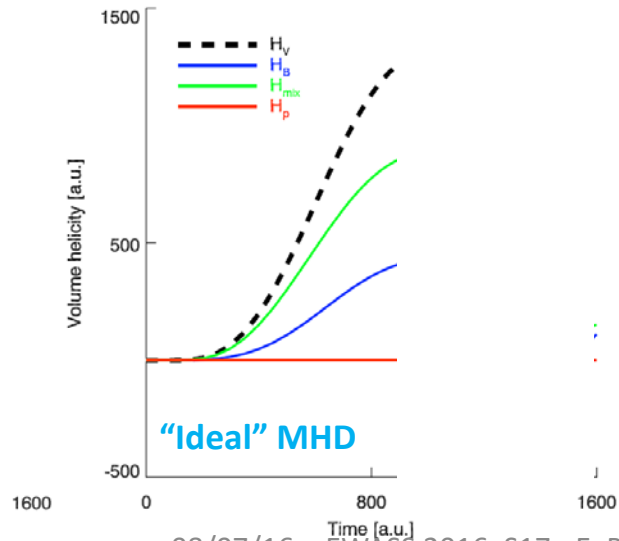
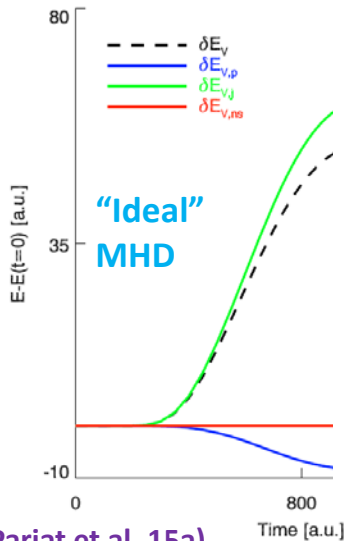
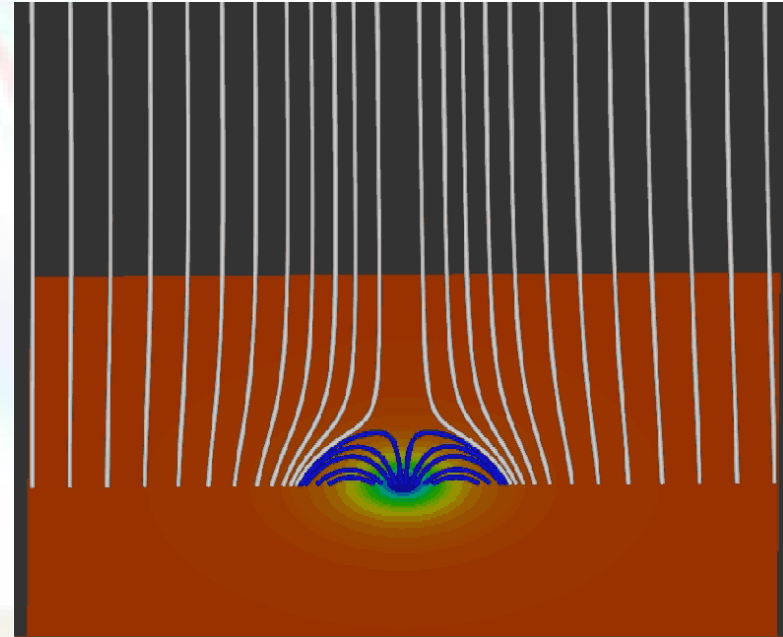
- **Helicity-conservation estimation:** measure the difference between
 - helicity variations in \mathcal{V}
 - helicity flux through the boundary sides \mathcal{S} .
- Method independent of the non-ideal processes, i.e. reconnection-model

$$C_m = -2 \int_{\mathcal{V}} \mathbf{E} \cdot \mathbf{B} \, d\mathcal{V} = \frac{dH}{dt} - F_{tot}$$

Test Case

(Pariat et al. 10)

- **3D MHD simulation of a solar coronal jet:**
Pariat et al. 09,10,15b ; Dalmasse et al. 12
 - Magnetic helicity/energy injected by bottom boundary motions
- **First phase: helicity/energy storage.**
 - Quasi-ideal MHD: reconnection inhibited.
- **Second phase: Jet generation**
 - Very impulsive energy release by recon.
 - Ejection of helicity.



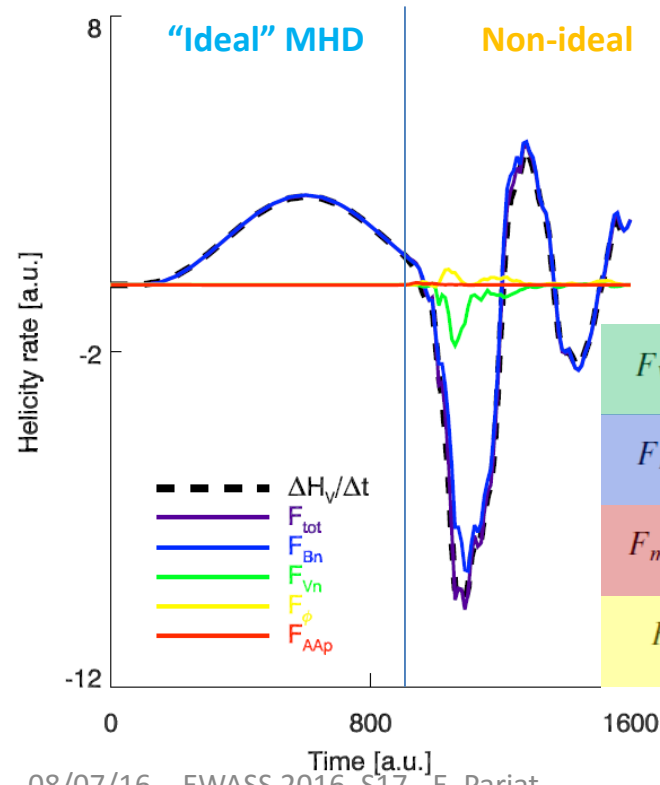
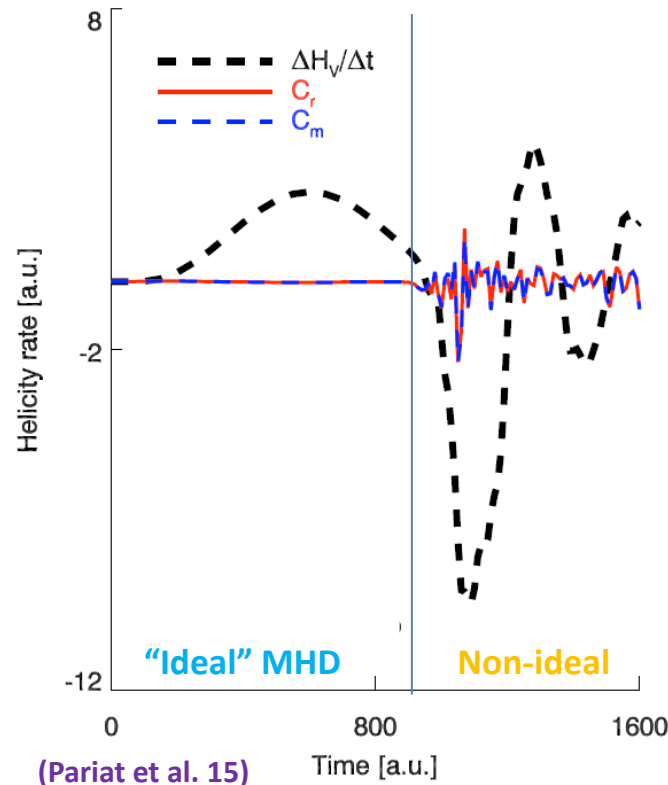
Helicity conservation

- Helicity and its flux are estimated independently
 - Direct volume helicity computation (Valori et al. 12): \mathbf{B} in \mathcal{V}
 - Helicity flux computation: \mathbf{B} , \mathbf{v} on \mathcal{S}
- ➔ Magnetic helicity is very well conserved both during the quasi-ideal MHD and non-ideal phases.

$$H_{\mathcal{V}} = \int_{\mathcal{V}} (\mathbf{A} + \mathbf{A}_p) \cdot (\mathbf{B} - \mathbf{B}_p) d\mathcal{V}$$

$$F_{tot} = F_{Vn} + F_{Bn} + F_{mix} + F_{\phi}$$

$$C_m = \frac{dH}{dt} - F_{tot}$$



$$F_{Vn} = 2 \int_{\partial\mathcal{V}} (\mathbf{B} \cdot \mathbf{A}) \mathbf{v} \cdot d\mathbf{S}$$

$$F_{Bn} = -2 \int_{\partial\mathcal{V}} (\mathbf{v} \cdot \mathbf{A}) \mathbf{B} \cdot d\mathbf{S}$$

$$F_{mix} = \int_{\partial\mathcal{V}} \left((\mathbf{A} - \mathbf{A}_p) \times \frac{\partial(\mathbf{A} + \mathbf{A}_p)}{\partial t} \right) \cdot d\mathbf{S}$$

$$F_{\phi} = -2 \int_{\partial\mathcal{V}} \frac{\partial\phi}{\partial t} \mathbf{A}_p \cdot d\mathbf{S}$$

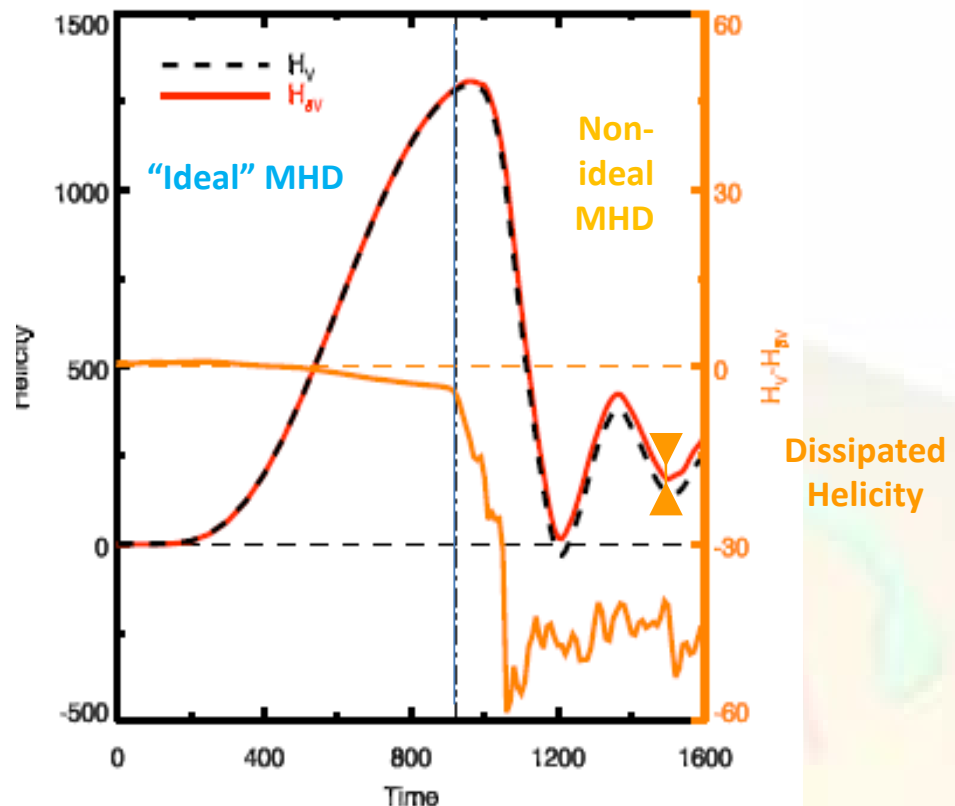
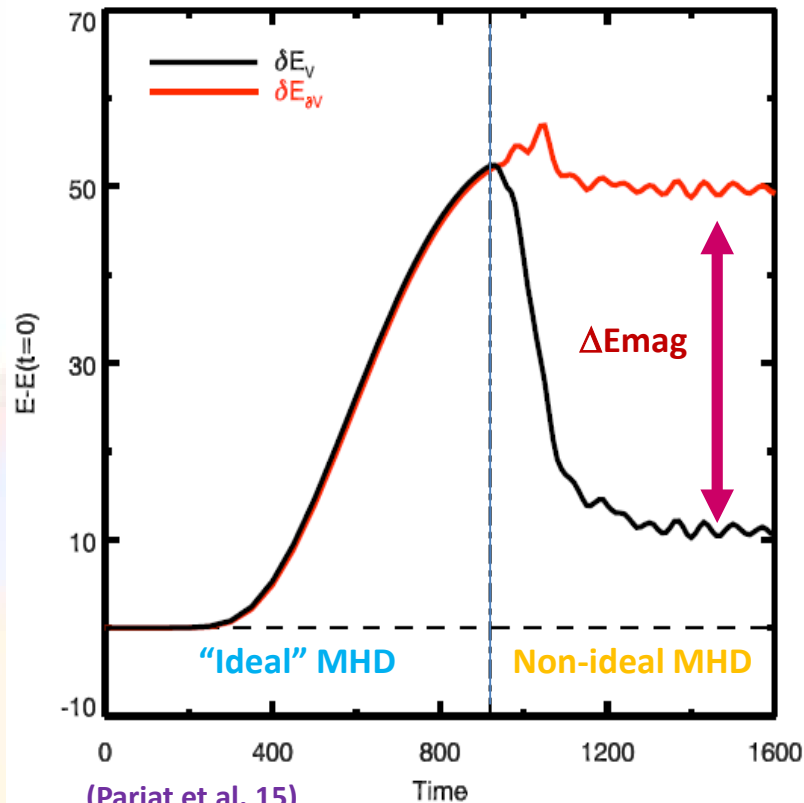
Helicity conservation

- **Magnetic helicity is very well conserved.**
 - Dissipated helicity is very small compared to the helicity injected in the system.
 - The dissipated helicity is very small compared to the amount of magnetic energy dissipated.

$$\frac{\Delta H_{\text{diss, Ideal}}}{\Delta H_{\text{inj}}} < 3 \times 10^{-3}$$

$$\frac{\Delta H_{\text{diss, Non-ideal}}}{\Delta H_{\text{ini}}} < 0.02$$

$$\frac{\Delta E_{\text{diss, Non-ideal}}}{\Delta E_{\text{inj}}} \sim 0.6$$



“Vn” and “Bn” flux terms

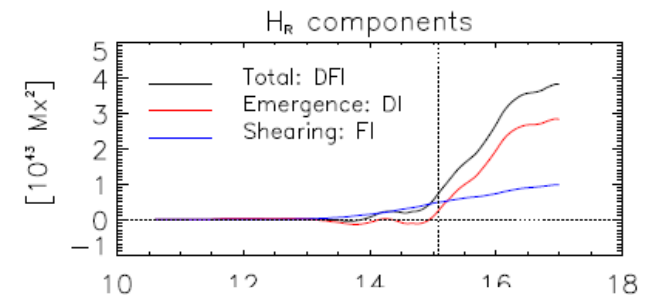
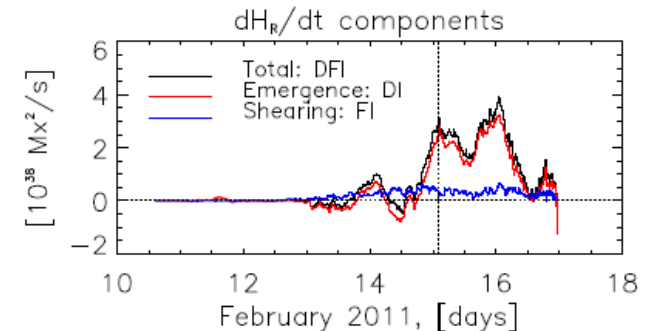
- The helicity flux decomposition presents two terms directly linked with the transport of helicity of the studied field
 - “emergence” term, depends on Vn
 - “shear” term, depends on Bn

$$F_{Vn} = 2 \int_{\partial V} (\mathbf{B} \cdot \mathbf{A}) \mathbf{v} \cdot d\mathbf{S}$$

$$F_{Bn} = -2 \int_{\partial V} (\mathbf{v} \cdot \mathbf{A}) \mathbf{B} \cdot d\mathbf{S}$$

$$F_{tot} = F_{Vn} + F_{Bn} + F_{mix} + F_{\phi}$$

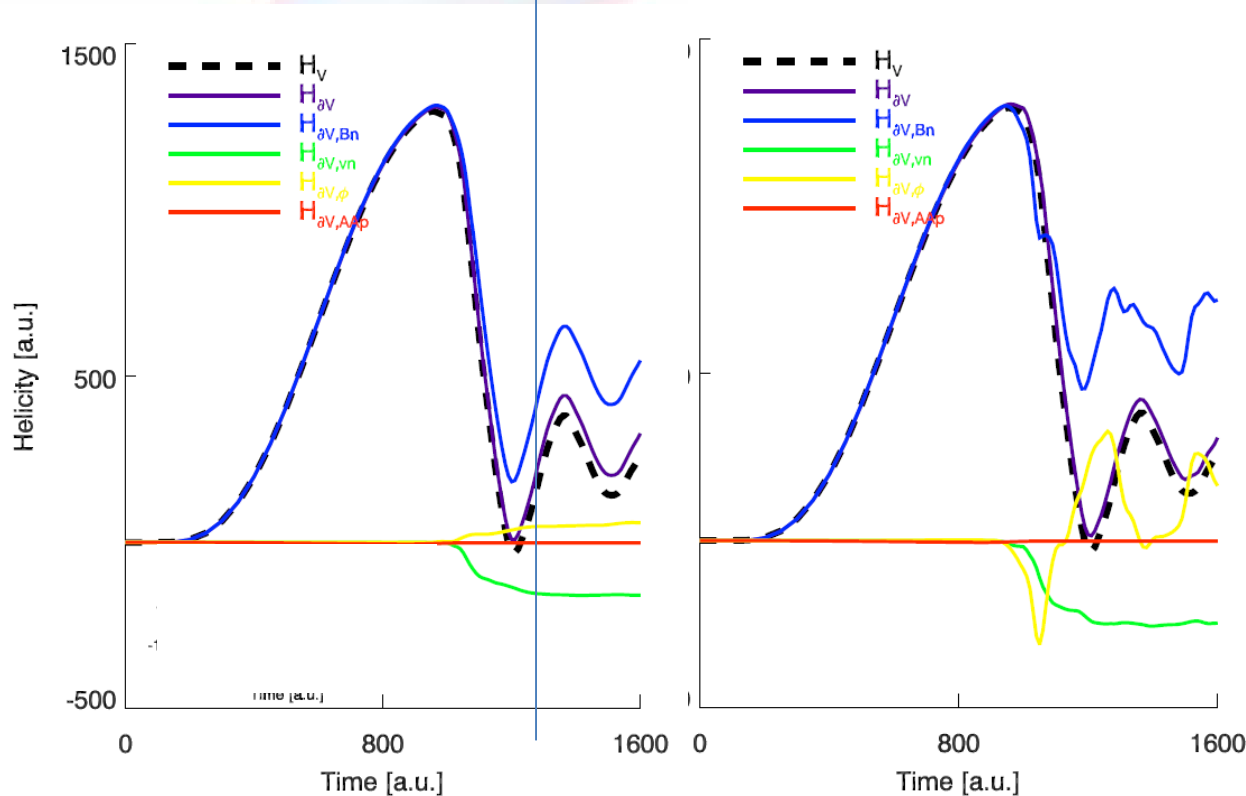
- Frequently used in the literature
 - Berger et al. 84, Maeshiro et al. 04, Pariat et al. 05, Démoulin & Pariat 09, Liu et al. 14a,14b; Kazachenko et al. 15, ...
- **1st problem: helicity is not a local quantity!**
 - Emergence term is different from self term of emerging structure
 - Emergence term is influenced by surrounding coronal field



Kazachenko et al. 15

“Vn” and “Bn” flux terms

- The “Vn” and “Bn” terms are not independently gauge invariant terms!
 - Their ratio can vary when computed with different gauge set (**A**, **A_p**)
 - Their sum also varies depending on the gauges
- ➔ the « Vn » and « Bn » terms do not convey any physical meaning



$$F_{Vn} = 2 \int_{\partial V} (\mathbf{B} \cdot \mathbf{A}) \mathbf{v} \cdot d\mathbf{S}$$

$$F_{Bn} = -2 \int_{\partial V} (\mathbf{v} \cdot \mathbf{A}) \mathbf{B} \cdot d\mathbf{S}$$

$$F_{mix} = \int_{\partial V} \left((\mathbf{A} - \mathbf{A}_p) \times \frac{\partial(\mathbf{A} + \mathbf{A}_p)}{\partial t} \right) \cdot d\mathbf{S}$$

$$F_{\phi} = -2 \int_{\partial V} \frac{\partial \phi}{\partial t} \mathbf{A}_p \cdot d\mathbf{S}$$

(Pariat et al. 15)

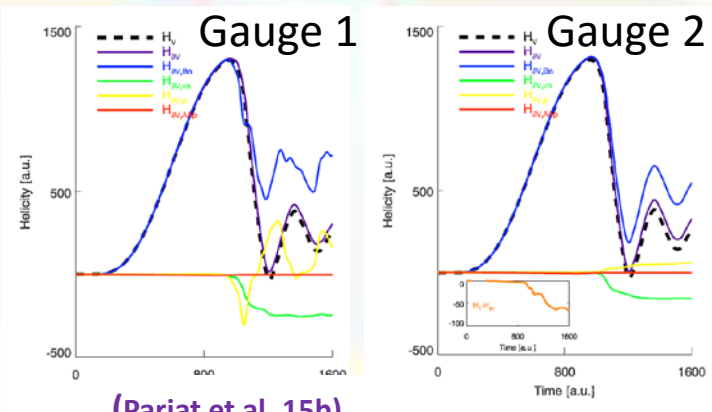
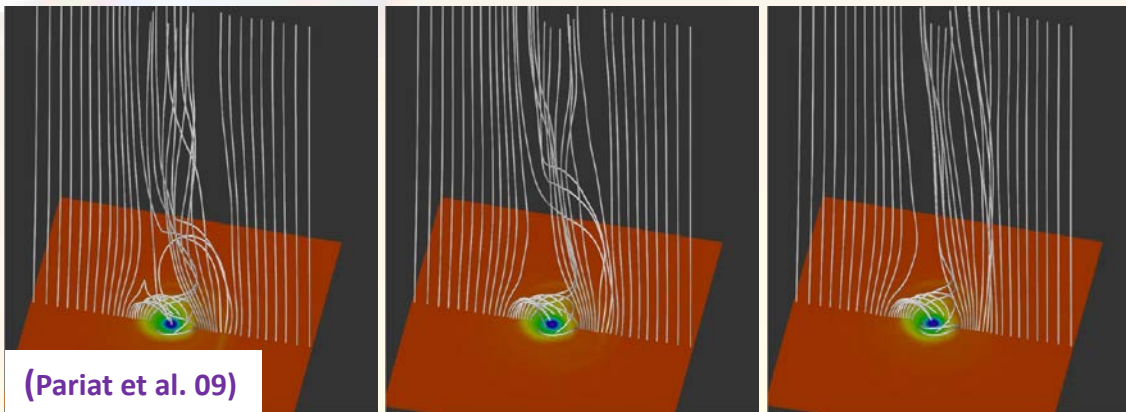
Gauges: **A**: DeVore ; **A_p**: DeVore + Coulomb

Gauges : **A** & **A_p**: DeVore

Conclusion



- Recent theoretical & numerical advances now allow the correct estimation of helicity in 3D numerical data sets.
- Estimations of the helicity conservation on an impulsive solar active like events (solar coronal jet).
 - Independent of reconnection models
 - Using several general gauges.
 - **As conjectured, magnetic helicity is very well conserved**
- **Forty years after, the Taylor conjecture can now be numerically tested in general configurations, using typical numerical data sets**
- ➔ Study and characterization of magnetic helicity in solar atmosphere: HELISOL





Thanks for your attention